EPJ Web of Conferences **81**, 02007 (2014) DOI: 10.1051/epjconf/20148102007

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Status of the measurement of $K_S \to \pi e \nu$ branching ratio and lepton charge asymmetry with the KLOE detector

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Abstract. We present the current status of the analysis of about 1.7 billion $K_S K_L$ pair events collected at DA Φ NE with the KLOE detector to determine the branching ratio of $K_S \to \pi e \nu$ decay and the lepton charge asymmetry. This sample is ~ 4 times larger in statistics than the one used in the previous most precise result, from KLOE as well, allowing us to improve the accuracy on the measurement and related tests of CPT symmetry and $\Delta S = \Delta Q$ rule.

1 Introduction

The \mathcal{CPT} symmetry assumes invariance of physical laws under the combination of the symmetries such as charge conjugation (C), parity (\mathcal{P}) and time reversal (\mathcal{T}) . One of possible ways to test violation of \mathcal{CPT} symmetry and basic assumptions of the Standard Model in the neutral kaon system is based on the difference between charge asymmetries for short-lived kaon (A_S) and for long-lived kaon (A_L) . Presently this difference is compatible with zero within errors, which suggests conservation of \mathcal{CPT} symmetry, however the value of A_L [1] was determined with a precision more than two orders of magnitude better than A_S [2].

2 Charge asymmetry and experimental verification

According to the Standard Model, weak force is responsible for semileptonic decay of K^0 or \bar{K}^0 . This implies that only two of four possible K^0 or \bar{K}^0 semileptonic decays occur (Figure 1, Table 1) and the change of strangeness (ΔS) entails the corresponding change of electric charge (ΔQ). This is the $\Delta S = \Delta Q$ rule. Semileptonic amplitudes can be parametrized as shown in Table 1 and connected to the conservation of discrete symmetries (Table 2) [3].

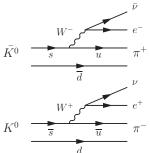


Figure 1. Feynman diagrams for K^0 and \bar{K}^0 semileptonic decays.

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Table 1. Relations between semileptonic decays of $K^0(\bar{K^0})$, introduced parametrization and discrete symmetries.

According to the Standard Model	Decay	Matrix element parametrization				
allowed	$K^0 \to \pi^- e^+ \bar{\nu}$	a+b	$= \langle \pi^- e^+ \nu H_{weak} K^0 \rangle$	$= (\text{if } \mathcal{CP}) = a^* + b^*,$		
allowed	$\bar{K^0} \rightarrow \pi^+ e^- \nu$	a^*-b^*	$= \langle \pi^+ e^- \bar{\nu} H_{weak} \bar{K^0} \rangle$	$= (\text{if } \mathcal{CP}) = a + b,$		
not allowed	$K^0 \rightarrow \pi^+ e^- \nu$	c+d	$= \langle \pi^+ e^- \bar{\nu} H_{weak} K^0 \rangle$	$= (if \mathcal{T}) = c^* - d^*,$		
not allowed	$\bar{K^0} \rightarrow \pi^- e^+ \bar{\nu}$	c^*-d^*	$= \langle \pi^- e^+ \nu H_{weak} \bar{K^0} \rangle$	$= (\text{if } \mathcal{T}) = c + d.$		

Table 2. Relations between discrete symmetries and semileptonic amplitudes.

		Conserved quantity					
		СР	\mathcal{T}	СРТ	$\Delta S = \Delta Q$		
er	a	Im = 0	Im = 0				
Parameter	b	Re = 0	Im = 0	= 0			
	c	Im = 0	Im = 0		= 0		
Pe	d	Re = 0	Im = 0	= 0	= 0		

Semileptonic amplitudes can be associated to the K_S and K_L semileptonic decay widths through the charge asymmetry:

$$\begin{split} A_{S,L} &= \frac{\Gamma(K_{S,L} \to \pi^- e^+ \nu) - \Gamma(K_{S,L} \to \pi^+ e^- \bar{\nu})}{\Gamma(K_{S,L} \to \pi^- e^+ \nu) + \Gamma(K_{S,L} \to \pi^+ e^- \bar{\nu})} \\ &= 2 \left[Re \left(\epsilon_K \right) \pm Re \left(\delta_K \right) + Re \left(\frac{b}{a} \right) \mp Re \left(\frac{d^*}{a} \right) \right] \\ &\text{if } \Delta Q = \Delta S \\ &= 2 \left[Re \left(\epsilon_K \right) \pm Re \left(\delta_K \right) + Re \left(\frac{b}{a} \right) \right] \\ &\text{if } C\mathcal{PT} \text{ and } \Delta Q = \Delta S \\ &= 2 \left[Re \left(\epsilon_K \right) \right]. \end{split}$$

The charge asymmetry for K_L was precisely determined from the KTeV experiment at Fermilab [1]:

$$A_L = (3.322 \pm 0.058_{stat} \pm 0.047_{syst}) \times 10^{-3},\tag{1}$$

while the most precise measurement of A_S was conducted by the KLOE collaboration [2]:

$$A_S = (1.5 \pm 9.6_{stat} \pm 2.9_{syst}) \times 10^{-3}.$$
 (2)

The obtained charge asymmetry for K_S decays is consistent, within error, with the charge asymmetry for K_L decays, which suggests conservation of \mathcal{CPT} symmetry. This result is dominated by the statistical uncertainty which is three times larger than the systematic contribution.

3 Measurement

The KLOE experiment is located at DAΦNE e^+e^- collider that works at the center of mass energy of the ϕ -meson mass ($\sqrt{s} = m_{\phi}$). The KLOE detector was optimized for efficient detection of long-lived kaons. A 2 m radius drift chamber allows to register around 40% of long-lived kaon decays inside the chamber while the rest reach the electromagnetic calorimeter. Identification of events with long-lived kaon ensures occurrence of short-lived kaon near the interaction point and vice versa. In order to improve signal over background ratio kinematic selection is applied. On remaining events the time-of-flight technique, which aims at rejecting background and identifying the final charge states ($\pi^+e^-\bar{\nu}$ and $\pi^-e^+\nu$), is used. Altogether about $10^5~K_S \to \pi e \nu$ decays were reconstructed, which will be used for the measurement of the charge asymmetry and branching ratio for K_S semileptonic decays. The analysis is still in progress, nevertheless it shows potential of reaching a twice better statistical error determination based on four times larger data sample. Also, due to the upgrade of KLOE detector and DAΦNE collider, further reduction of systematical and statistical uncertainties are expected in the future [4].

Acknowledgements

We warmly thank our former KLOE colleagues for the access to the data collected during the KLOE data taking campaign. We thank the DAΦNE team for their efforts in maintaining low background running conditions and their collaboration during all data taking. We want to thank our technical staff: G.F. Fortugno and F. Sborzacchi for their dedication in ensuring efficient operation of the KLOE computing facilities; M. Anelli for his continuous attention to the gas system and detector safety; A. Balla, M. Gatta, G. Corradi and G. Papalino for electronics maintenance; M. Santoni, G. Paoluzzi and R. Rosellini for general detector support; C. Piscitelli for his help during major maintenance periods. This work was supported in part by the EU Integrated Infrastructure Initiative Hadron Physics Project under contract number RII3-CT- 2004-506078; by the European Commission under the 7th Framework Programme through the 'Research Infrastructures' action of the 'Capacities' Programme, Call: FP7-INFRASTRUCTURES-2008-1, Grant Agreement No. 227431; by the Polish National Science Centre through the Grants No. DEC-2011/03/N/ST2/02641, 2011/01/D/ST2/00748, 2011/03/N/ST2/02652, 2013/08/M/ST2/00323, and by the Foundation for Polish Science through the MPD programme and the project HOMING PLUS BIS/2011-4/3.

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